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3	Retrieval of sun-induced fluorescence using advanced spectral fitting methods
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25 Abstract

The FLuorescence EXplorer (FLEX) satellite mission, candidate of ESA's 8th Earth Explorer program, 26 is explicitly optimized for detecting the sun-induced fluorescence emitted by plants. It will allow 27 28 consistent measurements around the O₂-B (687 nm) and O₂-A (760 nm) bands, related to the red and 29 far-red fluorescence emission peaks respectively, the photochemical reflectance index, and the 30 structural-chemical state variables of the canopy. The sun-induced fluorescence signal, overlapped to 31 the surface reflected radiance, can be accurately retrieved by employing the powerful spectral fitting 32 technique. In this framework, a set of fluorescence retrieval algorithms optimized for FLEX are 33 proposed in this study. Two main retrieval approaches were investigated: i) the optimization of the 34 spectral fitting for retrieving fluorescence at the oxygen absorption bands; ii) the extension of the 35 spectral fitting to a broader spectral window to retrieve the full fluorescence spectrum in the range from 36 670 to 780 nm. The accuracy of the retrieval algorithms is assessed by employing atmosphere-surface 37 radiative transfer simulations obtained by coupling SCOPE and MODTRAN5 codes. The simulated 38 dataset considers more realistic conditions because it includes directional effects, and the top-of-39 atmosphere radiance spectra are resampled to the current specifications of the *FLuORescence Imaging* 40 Spectrometer (FLORIS) planned to serve as the primary instrument aboard FLEX. The retrieval 41 accuracy obtained at the O₂-A band is strongly affected by directional effects, and better performance is 42 found in cases where directional effects are lower. However, the best performing algorithms tested provided similar performance, the RMSE (RRMSE) is 0.044 mW m⁻² sr⁻¹ nm⁻¹ (6.2%) at the O₂-A band, 43 0.018 mW m⁻² sr⁻¹ nm⁻¹ (2.9%) at the O₂-B band, and 6.225 mW m⁻² sr⁻¹ (6.4%) for the spectrally 44 integrated fluorescence emission. The promising results achieved open new perspectives extending 45 46 fluorescence studies not only in limited absorption bands, but its spectral behavior in relation to 47 different plant species, photosynthetic rates and stress occurrences.

48 **1** Introduction

49 Satellite remote sensing provides fundamental data to study and monitor vegetation state variables and 50 processes. Most Earth Observation missions rely on the analysis of the reflected radiance in the solar 51 domain to derive bio-physical (i.e. fractional cover (FC), leaf area index (LAI)), and bio-chemical 52 constituents of vegetation (i.e. chlorophyll, water and nitrogen). In the last few years, the remote 53 sensing of sun-induced fluorescence (SIF) represents a novel approach to provide new insight into plant 54 photosynthetic activity. SIF is a faint light signal released by the photosynthetic apparatus that plant 55 canopies add continuously to the reflected radiance in the visible and near-infrared wavelength range. 56 The strong interest and the various efforts ongoing by scientific community are prompted by the close 57 link between SIF and the actual plant photosynthetic rate (Baker, 2008; Papageorgiou & Govindjee, 58 2004). The exploitation of this signal at continental and global scale for examining the atmosphere-59 vegetation carbon exchanges estimations (Guanter et al., 2014; Lee et al., 2013; Parazoo et al., 2014), 60 and its parametrization into the Community Land Models (Lee et al., 2015), represent recent 61 applications.

62 In the last few years, global scale maps of SIF in the far-red region have been generated by exploiting 63 high spectral resolution sensors on board current space-borne mission, primarily devoted to atmospheric 64 chemistry. Joiner et al., (2011), Frankenberg et al., (2011) and Guanter et al., (2012) produced the first 65 global maps of far-red SIF by exploiting the data produced by the TANSO Fourier Transform 66 Spectrometer on board the Japanese GOSAT satellite (Hamazaki et al., 2005; Kuze et al., 2009). 67 Afterwards, global maps of far-red fluorescence have also been produced by exploiting the SCanning 68 Imaging Absorption SpectroMeter for Atmospheric CHartographY (SCIAMACHY) aboard of 69 ENVISAT (Joiner et al., 2012), and the Global Ozone Monitoring Experiment 2 (GOME-2) flying on 70 the operational European meteorological (MetOp) satellites (Joiner et al., 2013). Future advancements 71 for the far-red SIF region are expected from NASA's OCO-2 (Frankenberg et al., 2014) launched in

72 July 2014, and the forthcoming TROPOspheric Monitoring Instrument (TROPOMI) to be aboard the 73 Sentinel-5 Precursor (Guanter et al., 2015). The SIF retrieval algorithms from atmospheric space-borne 74 satellites are mostly based on the analysis of the absolute in-filling of fluorescence in the solar 75 Fraunhofer lines (740-755 nm). Even though the relative contribution of SIF is lower in such narrow 76 lines compared to broader and deeper telluric oxygen bands, the main advantage of this approach relies 77 in a simplified radiative transfer modeling of the atmospheric scattering and absorption effects in such 78 spectral windows. The SIF global maps derived from the atmospheric sensors can be very useful for a 79 global assessment of SIF, but their coarse spatial resolution (between few up to tens kilometers) do not 80 provide optimal data to study terrestrial ecosystems. In fact, the heterogeneity of the natural land 81 surface cannot be properly represented at these spatial resolutions in most of the cases. Furthermore, the 82 spectral configurations of the existing atmospheric missions (mostly between 757-775 nm) do not allow 83 the retrieval of red SIF, thus limiting the analysis to the far-red SIF only.

The *FLuorescence EXplorer* (FLEX) mission, candidate to the 8th Earth Explorer program, is currently 84 85 under Phase A/B1 study by the European Space Agency (ESA). FLEX is explicitly optimized for 86 detecting the SIF emitted by plants at a unique spatial resolution of 300 m (ESA, 2015). It will fly in 87 tandem with the ESA's Sentinel-3 (S3) (Donlon et al., 2012) to take advantage of complementary 88 measurements from the Ocean and Land Color Instrument (OLCI) and the Sea and Land Surface 89 Temperature Radiometer (SLSTR). The FLEX/S3 tandem mission will provide numerous advantages 90 including an accurate atmospheric correction, the consistent measurements of both red and far-red 91 fluorescence peaks, the detection of the photochemical reflectance index (PRI) (Gamon et al., 1992), 92 and estimation of bio-physical and bio-chemical canopy parameters. All these sources of information 93 are essential in this mission for a better understanding and interpretation of sun-induced fluorescence, 94 canopy variables and their relationships. Moreover, the canopy temperature delivered by S3 in combination with these information can allow to parametrize photosynthesis models to derive higherlevel products.

97 The FLuORescence Imaging Spectrometer (FLORIS) to fly aboard the FLEX satellite consists of two 98 spectrometers explicitly designed to provide systematic high-resolution spectral radiance observations 99 (0.3 nm) around the O₂-A (760 nm) and O₂-B (687 nm) absorption bands together with the continuous 100 spectral coverage in the visible to near-infrared spectrum (500-780 nm, ≤ 3 nm). The technical 101 specifications of FLORIS, in terms of spectral coverage, spectral resolution (SR), spectral sampling 102 interval (SSI) and signal to noise ratio (SNR), are well suited for retrieving fluorescence by using 103 spectral fitting methods (Mazzoni et al., 2012; Mazzoni et al., 2008; Mazzoni et al., 2010; Meroni et al., 104 2010). These methods make use of proper mathematical functions to simultaneously model the surface 105 reflectance and fluorescence at different wavelengths within spectral windows confined to the oxygen 106 absorption bands. However, FLORIS not only permits retrieving SIF at the two narrow O₂ bands with 107 high accuracy, but it also offers a continuous spectral coverage over the entire spectral region where the 108 fluorescence emission occurs. Taking advantage of this, spectral fitting approach can be further 109 improved and tested over broader spectral windows with the aim of recovering the entire fluorescence 110 emission spectrum. This could represent a relevant advancement because sun-induced fluorescence 111 spectrum is a complex function which mainly depends on the specific emission of photosystem I (PS_I) 112 and photosystem II (PS_{II}), and the successive re-absorption/scattering effects that occur at both leaf and 113 canopy levels. The variability of SIF spectrum at leaf level in relation to different plant species, 114 environmental conditions and plant's stress occurrence are documented in Agati, (1998) and Van 115 Wittenberghe et al., (2013). Therefore, the possibility of retrieving the fluorescence spectrum opens 116 novel and promising perspectives to have a better understanding of fluorescence in relation to 117 photosynthesis and other canopy state variables.

In this framework, the main aim of this study is to develop and testing SIF retrieval algorithms 118 119 optimized for FLORIS, and suitable for similar high-resolution sensors, based on Spectral Fitting 120 Methods. The specific objectives consist in: i) optimizing existing SFMs in confined spectral windows around the O₂ absorption bands; and ii) developing a novel retrieval algorithm which allows estimating 121 122 SIF spectrum. The paper is structured as follows: section 2 describes the FLEX mission providing 123 technical details on the FLORIS instrument. Section 3 describes the radiative transfer equations used, 124 the SIF retrieval algorithms developed, and the metrics used for evaluating the retrieval accuracy. 125 Section 4 shows and discuss the results with emphasis on the impact of atmosphere-surface directional 126 effects. Section 5 describes a number of indices that can be derived from FLEX. The main findings of 127 the study are summarized in section 6.

128

129 2 The FLEX/S3 tandem mission

130 The FLEX satellite is expected to fly in tandem with the Sentinel-3 in a Sun-synchronous orbit at an 131 altitude of about 815 km to deliver imagery at 300 m spatial resolution with a swath of 150 km (ESA, 132 2015). The revisit time will be 27 days at the Equator and more frequent acquisitions (\sim 19 days) over 133 high latitudes due to orbital overlaps. The lifetime foreseen of the mission is 3.5 years. The FLuORescence Imaging Spectrometer (Kraft et al., 2012, 2013), on board of FLEX, is a pushbroom 134 135 imaging spectrometer designed to detect canopy fluorescence and reflectance within a spectral range 136 between 500 and 780 nm. The current configuration of the FLORIS instrument consists of two 137 spectrometers: i) the Narrow Band Spectrometer (NBS) which provides high-resolution radiance in 138 defined spectral ranges around the O₂-A and O₂-B absorption bands, ii) the Wide Band Spectrometer 139 (WBS) characterized by a broader spectral coverage from 500 to 740 nm, with a lower spectral 140 resolution. The spectral bands are then resampled according to a defined binning scheme providing the 141 proper SR, SSI and SNR over the different spectral ranges. The resulting characteristics expected for a 142 typical spectral radiance observation by FLORIS are reported in Table 1. The required SNR values 143 synthetized in the table refer to the spectral bands, therefore the expected actual values will be higher 144 after the spectral binning.

145

146 <Table 1>

147

The combination of NBS and WBS spectra will provide both the high SR for fluorescence retrieval at the O_2 bands, and a broader spectral coverage, at the same time (Figure 1). This particular combination of variable SRs, SSI, and the binning-scheme will provide an unprecedented SNR, which is one key factor ensuring accurate retrieval of SIF.

152

153 <Figure 1>

154

155 **3** Materials and Methods

156 **3.1 Radiative transfer simulations**

157 A dataset of radiative transfer (RT) simulations has been initially created to develop the retrieval 158 algorithms and to assess their performances. It consists of fluorescence, reflectance and total upward 159 radiance at top of atmosphere (TOA) and bottom of atmosphere (BOA, assumed to be equivalent to 160 TOC, top-of-canopy) calculated by coupling the Soil Canopy Observation Photosynthesis Energy 161 balance (SCOPE) (Van Der Tol et al., 2009) with the MODerate resolution atmospheric TRANsmission (MODTRAN) RT models. In particular, the forward model used relies on the four-stream radiative 162 163 transfer theory (Verhoef & Bach, 2007, 2012) with the addition of the direct and diffuse fluorescence 164 fluxes (Verhoef et al., 2014). It represents an accurate, efficient and relatively simple way to describe 165 the radiative transfer interactions of the Earth's surface-atmosphere system. The surface-atmosphere 166 RT interactions and the subsequent propagation to the top of the atmosphere (L^{TOA}) are defined in the 167 following radiative transfer equation, Eq. 1.

168

$$L^{\text{TOA}} = \rho_{so} \frac{E_s^0 \cos \theta_s}{\pi} + \left[\frac{\tau_{ss} r_{so} E_s^0 \cos \theta_s}{\pi} + SIF_s + \frac{(\tau_{sd} + \tau_{ss} \overline{r_{sd}} \rho_{dd}) E_s^0 \cos \theta_s / \pi + \overline{SIF_d} \rho_{dd}}{1 - \overline{r_{dd}} \rho_{dd}} r_{do} \right] \tau_{oo}$$
Eq. 1
$$+ \left[\frac{(\tau_{sd} \overline{r_{dd}} + \tau_{ss} \overline{r_{sd}}) E_s^0 \cos \theta_s / \pi + \overline{SIF_d}}{1 - \overline{r_{dd}} \rho_{dd}} \right] \tau_{do}$$

169

 L^{TOA} is therefore composed of three additive terms (the 3 lines on the right-hand side of Eq. 1) which 170 171 are referred to as the atmospheric path radiance, the target's surface radiance and the adjacency effect, respectively. The surface reflectance is modelled by four terms: r_{so} is the target bi-directional 172 reflectance factor (BRF), r_{do} the target hemispheric-directional reflectance factor (HDRF), $\overline{r_{sd}}$ the 173 spatially filtered directional-hemispherical reflectance factor (DHRF) of the surrounding, $\overline{r_{dd}}$ the 174 175 spatially filtered bi-hemispherical reflectance (BHRF) of the surrounding (c.f. Nicodemus et al., (1977) for a description of used reflectance quantities). The term ρ_{so} is the atmospheric bi-directional 176 177 reflectance and ρ_{dd} is the spherical albedo at the bottom of the atmosphere. The term τ_{ss} is the direct 178 atmospheric transmittance in the direction of the sun, τ_{oo} is the direct atmospheric transmittance in the 179 direction of viewing, τ_{sd} the diffuse atmospheric transmittance for solar incidence, and τ_{do} the directional atmospheric transmittance for diffuse incidence. The quantity E_s^0 is the extra-terrestrial solar 180 spectral irradiance on a plane perpendicular to the sunrays, and θ_s is the local solar zenith angle (SZA). 181 SIF_s is the sun-induced fluorescence radiance of the target in the observer's direction and πSIF_d the 182 183 hemispherical fluorescence flux of the surroundings. The RT simulations consider a homogenous 1-D 184 target in the horizontal plane, and the over bar indicates the spatial filtering of the terms related to the 185 infinitely extended surrounding area.

Strictly, Eq. 1 is only valid for monochromatic radiation because the atmospheric quantities ρ_{so} , ρ_{dd} , 186 τ_{ss} , τ_{oo} , τ_{sd} , and τ_{do} are strongly modulated by narrow absorption lines due to various gases in the 187 atmosphere, and therefore they are strongly correlated in spectral regions where many absorption lines 188 189 are present. Consequently, the mean value of a product of these quantities over a finite spectral interval 190 (i.e. integration of radiance in the sensor spectral bands) is not equal to the product of the mean values of the individual quantities. Hence, the products of the different atmospheric quantities needed are 191 192 formed and stored at high spectral resolution (Verhoef et al., 2014), before the convolution to the sensor 193 spectral band. The set of atmospheric transfer functions required by Eq. 1 is summarized in Figure 2.

194

195 <Figure 2>

196

197 Based on these transfer functions, Eq. 1 can be rewritten as (Eq. 2).

198

$$L^{TOA} = t_1 \left[t_2 + t_8 r_{so} + \frac{t_9 + t_{14} \overline{r_{sd}}}{1 - \overline{r_{dd}} t_3} r_{do} + \frac{t_{10} \overline{r_{sd}} + t_{11} \overline{r_{dd}}}{1 - \overline{r_{dd}} t_3} \right] + t_6 SIF_s + \left[\frac{SIF_d (t_7 + t_{13} r_{do})}{1 - \overline{r_{dd}} t_3} \right]$$
Eq. 2

199

200 The corresponding upward radiance at the bottom of atmosphere L^{BOA} is calculated using a scheme in 201 which some transfer functions represent single-way (t_4 and t_5) instead of two-way transmittances: 202

$$L^{BOA} = t_1 \left[t_4 r_{so} + \frac{t_5 + t_{12} \overline{r_{sd}}}{1 - \overline{r_{dd}} t_3} r_{do} \right] + SIF_s + \left[\frac{\overline{SIF_d} t_3 r_{do}}{1 - \overline{r_{dd}} t_3} \right]$$
Eq. 3

Eq. 3 can also be used to calculate the upwelling radiance from the White Lambertian Reference (L^{WLR}) panel by setting the reflectance terms from the target r_{so} and r_{do} equal to one and the target fluorescence equal to zero. This quantity represents the global (direct and diffuse) incoming radiance at ground level including the adjacency effects from the surrounding.

208

$$L^{WLR} = t_1 \left[t_4 + \frac{t_5 + t_{12}\overline{r_{sd}}}{1 - \overline{r_{dd}} t_3} \right] + \left[\frac{\overline{SIF_d} t_3}{1 - \overline{r_{dd}} t_3} \right]$$
Eq. 4

209

210 The ratio of L^{BOA} and L^{WLR} , as measured by ground-based measurements, is the surface apparent 211 reflectance (ρ_{app}), which includes both reflected and fluorescence radiance (Eq. 5). The surface 212 reflectance free of the fluorescence contribution (ρ), (Eq. 6), represents the reference value of the 213 retrieval algorithms.

214

$$\rho_{app} = \frac{L^{BOA}}{L^{WLR}}$$
 Eq. 5

$$\rho = \frac{L^{BOA} - SIF_s}{L^{WLR}}$$
 Eq. 6

215

The four surface reflectance terms, SIF_s and SIF_d spectra were simulated by using version 1.40 of the SCOPE model. The RT calculations in the solar-reflective domain are performed in SCOPE with a spectral sampling and resolution of 1 nm over the 400–2400 nm range. The atmospheric transfer functions were instead derived from the MODTRTAN 5.2.1 code (Berk et al., 2011). Different atmospheric conditions (i.e., visibility, humidity, aerosol type, profile etc.) were simulated at 1 cm⁻¹ spectral sampling over the 400–50000 nm range, and the atmospheric transfer functions were calculated by using the MODTRAN Interrogation Technique (MIT) (Verhoef & Bach, 2012). 223 In summary, the database consists of 31 cases simulated with different soil, leaf, canopy and 224 atmospheric parameters (Table 2). The values reported in the column labeled "medium" are used to 225 simulate case n° 19, which refers to a typical scenario in terms of vegetation and atmospheric 226 parameters. The other cases are deviations from the typical scenario and they are obtained by changing 227 mostly one parameter at a time. In particular, the cases 1-18 are deviations of surface parameters, while 228 cases 20-31 correspond to deviations of the atmospheric conditions, including SZA. The MODTRAN 229 simulations were done independently, but the SCOPE parameters were set-up consistently with 230 MODTRAN at corresponding values (indicated in Table 2 in the column 'Coupling'). It should be 231 noted that this novel RT dataset includes a fundamental advancement, since for the first time, it includes 232 a realistic modelling of the canopy-atmosphere directional effects by considering the four reflectance 233 terms and their coupling with the atmospheric functions. Nadir viewing was used for all the 234 simulations, which is close to reality for the satellite observations. The resulting spectra were convolved 235 with the FLORIS-NBS and FLORIS-WBS instruments spectral response functions, and they were 236 binned according to the current specifications provided by ESA. A Gaussian distributed noise, of which the variance was linearly related to the L^{TOA} intensity at the different wavelengths, was added to the 237 238 TOA radiances (Verhoef et al., 2014) to simulate the noise levels expected for FLORIS. The 239 variabilities of surface ρ , SIF and the TOA radiance simulated in the FLORIS configuration are 240 depicted in Figure 3.

The FLEX/S3 tandem mission will allow the exploitation of the additional data available from Sentinel-3 (e.g., in the blue and the cirrus bands) to provide a highly accurate atmospheric correction. The mission configuration is well suited to consider a two-step retrieval approach which considers a preliminary correction of the atmospheric effects followed by the decoupling of the SIF and reflectance. The two successive steps are executed independently starting from the TOA radiance to calculate the canopy fluorescence and reflectance. Specifically, the TOA is converted to the BOA radiance by the 247 atmospheric correction calculations, and then the fluorescence and reflected radiance are decoupled 248 from the BOA spectra. This approach has the advantage of a limited number of model parameters to be 249 estimated in the two distinct retrieval processes, thus providing stable results. In this way, the sources 250 and the magnitudes of the retrieval error for the atmospheric correction and the following fluorescence 251 retrieval can be better quantified and understood. It must be noted that the retrieval of the atmospheric 252 parameters, which includes a detailed coupling between the FLEX and S3 data, is not considered in this 253 work. A slightly simplified, but still relatively accurate, forward propagation model (Eq. 7) for the TOA 254 radiance is considered in this study as baseline for simulating the correction of atmospheric effects 255 (Verhoef et al., 2014). The right-hand side of Eq. 7 contains a linear combination of the four reflectance 256 terms and the two fluorescence terms in the numerator.

257

$$L^{TOA} \approx \left[t_1 t_2 + \frac{t_1 \left(t_8 r_{so} + t_9 r_{do} + t_{10} \overline{r_{sd}} + t_{11} \overline{r_{dd}} \right) + t_6 SIF_s + t_7 SIF_d}{1 - t_3 \overline{r_{dd}}} \right]$$
Eq. 7

258

259 The atmospheric correction applied is based on the assumption of a uniform and Lambertian surface 260 with a reflectance ρ and a fluorescent radiance SIF. The atmospheric correction is applied considering 261 the following quantities:

$$L_0 = t_1 t_2 Eq. 8$$

$$g\rho = t_1 \left(t_8 r_{so} + t_9 r_{do} + t_{10} \overline{r_{sd}} + t_{11} \overline{r_{dd}} \right)$$
 Eq. 9

$$tSIF = t_6 SIF_s + t_7 SIF_d$$
 Eq. 10

$$S = t_3$$
 Eq. 11

where the path radiance for a black surface with zero albedo, L_0 , is given by the t_1t_2 product (Eq. 8), gis the gain factor to obtain TOA radiance, ρ is the "effective" reflectance, i.e. a weighted average of the four reflectance terms (Eq. 9), tF is the "effective" transmitted fluorescence (Eq. 10), and S is the atmospheric spherical albedo (Eq. 11). With these quantities, the simplified atmospheric forward modelcan be rewritten as:

267

$$L^{TOA} \approx L_0 + \frac{g\rho + tSIF}{1 - \rho S}$$
 Eq. 12

268

269 The surface apparent reflectance R_{ac} (i.e. reflectance and SIF contributions) after atmospheric 270 correction (under the assumption of zero *SIF*) can be obtained by solving ρ from Eq. 12 for *SIF* = 0, and 271 turns out to be given by

272

$$R_{ac} = \frac{g \rho + t SIF}{g + S t SIF}$$
Eq. 13

273

The resulting R_{ac} is practically equal to a weighted average of the four surface reflectance factors, with a small contribution due to fluorescence from the target and the surrounding. The contribution from fluorescence will be relatively larger in atmospheric absorption bands, since *t* is less attenuated by absorption than *g*. After the atmospheric correction, considering a "perfect knowledge" of the atmospheric transfer functions, the best available approximation for the radiance of the white Lambertian reference is given by Eq. 14.

280

$$L^{WLR} = \frac{t_1(t_4 + t_5)}{1 - SR_{ac}}$$
 Eq. 14

281

Finally, the BOA radiance is obtained from Eq. 15 by combining Eq. 13 and Eq. 14. This means that L^{BOA} derives directly from L^{TOA} once the atmospheric parameters are known. In this way, the noise in L^{TOA} is also propagated into L^{BOA} .

285

$$L^{BOA} = R_{ac} L^{WLR}$$
 Eq. 15

It should be noted that the two-step procedure described above was applied in the numerical experiments discussed in this paper. However, on the basis of Eq. 12 one could also develop a single step approach, in which measured TOA radiances would be fitted using Eq. 12 with atmospheric transfer functions and modeled spectral curves of surface reflectance and fluorescence as inputs. In a single step approach no atmospheric correction has to be carried out, only atmospheric characterization, and the assumptions usually made for correction approaches (i.e. uniform and Lambertian surface, homogeneous atmosphere, etc.) are no longer necessary, leading to greater flexibility and robustness.

293 **3.2 Fluorescence retrieval**

Once the TOA radiance from FLORIS is converted to BOA after the atmospheric correction, the radiance spectra are now composed of the additive contributions of fluorescence and sun reflected radiance (Eq. 16).

297

$$L^{BOA}(\lambda) = \rho L^{WLR}(\lambda) + SIF(\lambda)$$
Eq 16

298

299 The decoupling of the two terms is then achieved by using the spectral fitting approach. Basically, it 300 consists in the optimization of the fit between modelled and measured spectra of both fluorescence and 301 reflectance by adjusting the mathematical functions used to describe SIF and ρ spectral behaviors 302 within the defined spectral window. The parameters of the mathematical functions used for SIF and p 303 are estimated through a least square nonlinear curve-fitting optimization technique that minimizes the 304 cost function in Eq. 17. The MATLAB function LSQCURVEFIT was used to perform the optimization. This routine allows to include upper and lower bounds for model parameters excluding solutions 305 306 without a physical meaning (i.e. negative SIF and ρ).

307

$$\min \sum \left(L^{BOA}(\lambda) - SIF(\lambda) - \rho(\lambda)L^{WLR}(\lambda) \right)^2$$
Eq. 17

308

309 The usage of a large number of spectral bands has many advantages allowing estimation of a larger 310 number of model parameters that describe the spectral variables behavior, and reduce the impact of 311 instrumental noise. A source of error in these SIF retrieval methods could be due to the wrong 312 assumptions (i.e. linear, polynomial) about spectral functions used in modelling SIF and ρ . In general, 313 the use of functions with a larger number of parameters increases the modelling capability of the 314 spectral variable of interest. It is usually required for modelling broader spectral windows, and it is 315 particularly relevant to account for sharp spectral features within the considered spectral range. On the 316 other hand, it could lead to undesired problems like data over-fitting (i.e., the fitting model does not 317 describe the variable of interest but the random error) or ill-posed numerical inversions. For these 318 reasons, an optimal balance between model parameters, spectral window, and retrieval accuracy is 319 needed and it is assessed in this study.

The spectral fitting method has been applied in the narrow spectral range provided by the FLORIS-NBS (hereafter named SFM according the terminology originally proposed in Mazzoni et al., (2012); Meroni et al., (2010)) to retrieve fluorescence in confined spectral windows centered at the O_2 -bands, while a novel spectral fitting based algorithm aimed at the retrieval of the full fluorescence emission spectrum and a series of derived products is hereafter referred to as *SpecFit*.

325

326 **3.2.1** Fluorescence retrieval at the O₂ bands (SFM)

327 The high resolution spectra around the two O_2 absorptions from the FLORIS-NBS instrument are 328 particularly indicated for the retrieval of SIF by exploiting spectral fitting methods (SFM). The 329 rationale behind the exploitation of narrow spectral regions, characterized by strong absorptions, resides 330 in the higher contribution of SIF with respect to the total radiance. For this reason, the SFM algorithms 331 are focused on the detection of SIF at the O₂ bands or, at least, in relatively narrow regions around the 332 main absorption features. The use of such fitting windows permits to reduce the impact of instrumental 333 noise (which is higher in the absorption bands), and to exploit additional absorptions features nearby 334 (i.e. the solar Fraunhofer lines in the 740-759 nm range). The modelling of SIF and ρ as a function of 335 wavelengths is simpler in relatively narrow spectral ranges, but less spectral information is used in the 336 decoupling. This is particularly true at the O_2 -A band where reflectance is a smooth function, while it is 337 more difficult at O₂-B due to the rising of the red-edge reflectance. In fact, the O₂-B band is very close 338 to the red-edge transition and it is located adjacent to the maximum of the fluorescence emission peak 339 (684-685 nm). Such behavior makes the modelling of the spectral functions difficult in this region. 340 Moreover, the typical radiance levels for vegetated surfaces in the red spectrum (i.e., chlorophyll 341 absorption) are generally much lower than in the near-infrared (around the O₂-A band) and 342 consequently the SNR is worse. On the other hand, the relative contribution of SIF to the total radiance 343 is larger in this spectral region. To evaluate all these aspects, two different spectral windows centered at 344 the main oxygen absorption bands (Table 3) were considered for evaluating these factors. In this 345 setting, narrow spectral windows within the FLORIS-NBS spectrum are exploited to retrieve SIF by 346 SFM.

347

348 < Table 3>

349

A total of 54 combinations of spectral functions were tested to predict fluorescence and reflectance (Table 4). The ρ models are based on polynomial functions (P), Legendre (L) polynomials, and piecewise cubic spline (S) with different knots (i.e., 2, 4, 6 knots). The SIF is represented with both polynomial (P), Legendre (L) and spline (S) functions (for models IDs 1-9) and by using Gaussian, Lorentzian and Voigt profiles (for model IDs 10-18). The models 1-9 for each class are obtained by combining functions with a different polynomial order (i.e., from 1st to 3rd degree) or number of knots for the cubic splines. The models 10-18 were obtained by considering polynomials for reflectance and Gauss, Lorentz and Voigt profiles to model fluorescence.

358

359 <Table 4>

360

361 **3.2.2** Retrieval of the full fluorescence spectrum (SpecFit)

With this approach we extend the spectral fitting technique over the entire spectral region where the 362 363 fluorescence emission occurs. It relies on the inversion of the radiance spectrum in the red to far-red 364 spectral region for estimating the full fluorescence spectrum in the 670-780 nm range. Also this approach takes advantage of the spectral configuration of FLORIS instrument. In particular, the high 365 resolution spectra around the O₂ bands from FLORIS-NBS are joined in the red-edge by the lower 366 367 resolution spectra from FLORIS-WBS in order to obtain continuous spectral coverage. The idea 368 underlying this algorithm is almost the same as the spectral fitting described earlier, but the 369 mathematical functions used to predict SIF and p spectral behaviours in such a broader spectral window 370 are more complex.

The piecewise cubic spline was selected to reproduce the reflectance signature in the red to far-red spectral region. The two red and far-red SIF emission peaks were modelled using different combinations of Gaussian, Lorentzian and Voigt profiles. Generally, the fitting of the red peak is simpler because almost all of the contribution is due to the PS_{II} only; PS_{I} only contributes to the rededge tail of the first (red) emission peak (Franck et al., 2002). This fact results in an almost stable position and width of the red peak. On the contrary, the modelling of the far-red peak is more difficult because there are both PS_{II} and PS_{I} contributions, and these two act differently. Depending on the 378 relative contribution of fluorescence emitted by PS_I or PS_{II}, the wavelength of the maximum emission 379 shifts over several nanometers (i.e., up to 5-6 nm) and in some cases the peak shows an asymmetric 380 behavior. Two different implementations were considered to achieve an accurate fitting of the full SIF 381 spectrum: i) the Lorentzian and Gaussian functions were used to model red (684 nm) and far-red (740 382 nm) peaks, respectively; ii) two Voigt functions with an asymmetry parameter were used for both 383 peaks. The first implementation method has the advantage of being more robust because a total of 6 384 parameters are used (i.e. 3 for each emission peak). However, it must be noted that Lorentzian and 385 Gaussian functions do not provide as much flexibility to obtain accuracy in modelling the peaks' 386 spectral profiles, as do the Voigt functions. Unfortunately, the computation of the Voigt function is time 387 expensive due to the convolution between the Gaussian and Lorentzian functions. This point is critical 388 when the functions must be computed several times within the iterative optimization process. For this 389 reason the so-called pseudo-Voigt function, computed as a weighted combination of the Lorentzian and Gaussian functions, was selected for modelling each one of the two fluorescence emission peaks. The 390 391 first term of the pseudo-Voigt function represents the Lorentzian contribution, while the second term 392 represents the Gaussian contribution. The full SIF emission spectrum is therefore the sum (Eq. 18) of 393 the two pseudo-Voigt functions. The pseudo-Voigt functions were used to model the red (Eq. 19) and 394 far-red (Eq. 20) peaks respectively. The far-red peak includes also an additional parameter (Eq. 21) to 395 account for the peak asymmetry as proposed in Stancik & Brauns, (2008).

$$SIF(\lambda) = SIF_{far-red}(\lambda) + SIF_{red}(\lambda)$$
 Eq. 18

$$SIF_{red}(\lambda) = f \frac{\mu}{\left(\frac{\lambda - \lambda_0}{\sigma(\lambda)}\right)^2 + 1} + (1 - f) \ \mu \exp\left(-\frac{(\lambda - \lambda_0)^2}{2\sigma(\lambda)^2}\right)$$
Eq. 19

$$SIF_{far-red}(\lambda) = f \frac{\mu}{\left(\frac{\lambda - \lambda_0}{\sigma(\lambda)}\right)^2 + 1} + (1 - f) \mu \exp\left(-\frac{(\lambda - \lambda_0)^2}{2\sigma_{asym}(\lambda)^2}\right)$$
Eq. 20

$$\sigma_{asym}(\lambda) = \frac{2\sigma}{(1 + \exp(a(\lambda - \lambda_0)))}$$
Eq. 21

3.3 Error estimation

A number of statistical indicators were used to evaluate the performance of the retrieval algorithms in a consistent way. The accuracy was assessed by comparing the reference simulated SIF and p against the retrieved values at different λ . The statistical indexes considered in this study are the Root Mean Square Error (RMSE) (Eq. 22) which quantifies the amount by which an estimation differs from the assumed true value of the quantity, and the Relative RMSE (RRMSE) that represents the percentage of error with respect to the actual values (Eq. 23). The absolute difference (SIF_{diff}) between simulated and retrieved fluorescence (Eq. 24) provides information on whether retrieved fluorescence is underestimated or overestimated. Finally, a comparison was conducted of the coefficient of determination (r^2) between simulated and retrieved values.

$$RMSE = \sqrt{\frac{\sum_{\lambda=1}^{n} (SIF(\lambda) - SIF_{ret}(\lambda))^{2}}{n}}$$
Eq. 22
$$RRMSE = \sqrt{\frac{\sum_{\lambda=1}^{n} \left(\frac{SIF(\lambda) - SIF_{ret}(\lambda)}{SIF(\lambda)}\right)^{2}}{n}}{n}}$$
Eq. 23
$$SIF_{diff} = \int_{\lambda} SIF_{ret}(\lambda)d\lambda - \int_{\lambda} SIF(\lambda)d\lambda$$
Eq. 24

410

411 **4 Results and Discussions**

412 **4.1** Fluorescence at the O₂ bands

Although different spectral ranges were tested (data not shown), the broader range indicated in Table 3 was the most efficient for the retrieval of fluorescence. This is probably due to the additional information content in the broader spectral window (i.e., inclusion of solar Fraunhofer lines), and at the same time, reducing the impact of instrumental noise by using a larger set of spectral channels. The overall accuracy for the 31 cases in shown in Figure 4.

418

419 <Figure 4>

420

The overall accuracy depends on the performances of the spectral functions used for representing SIF and ρ in the numerical inversion. As observed in previous studies, the reflectance around the O₂-A band is generally smooth and most polynomials and Legendre polynomials of 2nd and 3rd order and piecewise cubic splines produce a suitable fit of the reflectance spectra (Mazzoni et al., 2012; Meroni et al., 2010). On the other hand, it can be observed that polynomial functions do not produce as accurate results in modelling the fluorescence spectrum when compared to peak-like functions (i.e., Gaussian, Lorentzian and Voigt). In fact, the SFM versions 1-9 resulted in larger RMSE values, with the only exception being 428 the piecewise cubic spline at the O_2 -A band. In summary, better accuracy was achieved by the SFM 429 version "18-Spline" which makes use of piecewise cubic spline and Voigt functions to model 430 reflectance and fluorescence respectively.

A different situation happened at the O_2 -B band, where sharp variations of both reflectance and fluorescence occur. The piecewise cubic spline and the Voigt spectral functions enabled SIF to be retrieved with a proper accuracy at the O_2 -B band using the SFM version "18-Spline". Since the same functions provided higher accuracy in both the O_2 -A and O_2 -B bands, the "18-Spline" should be considered as the best candidate for further implementation for FLEX. The further analysis hereafter refers to this specific SFM version.

437 The analysis of the retrieval performances for the different RT simulations is fundamental assessing the 438 sensitivity of the retrieval algorithm to specific parameters in the forward model. The retrieval 439 accuracies for each one of the different 31 RT simulations (Figure 5) are not significantly affected when 440 surface or atmospheric parameters are modified. For the O₂-A band, an average RMSE lower than 0.1 $mWm^{-2}sr^{-1}nm^{-1}$ (RRMSE% < 8) was found for almost all cases, with the exception of simulations 15-18 441 and 31 where significantly better performances were found with RMSEs < 0.025 mWm⁻²sr⁻¹nm⁻¹ 442 (RRMSEs% < 1). These cases represent extreme LAI values (0.5 and 6.0 m^2/m^2 for cases 15 and 16 443 444 respectively), leaf angle distribution (planophile and erectophile for cases 17, 18 respectively) and very 445 low solar zenith angles (case 31). The light penetration within the canopy, and the successive 446 interaction with soil, are key factors which cause most of the canopy directional effects. In fact, 447 low/high LAI values or erectophile/planophile leaves distributions determinate that light rays mostly 448 interact with soil/canopy only. On the contrary, intermediate cases show larger directional effects 449 caused by a larger mixing between the different soil/canopy contributions. The SIF_{diff} index shows that in the best cases the SIF values at O₂-A band are generally underestimated by about -0.1 mWm⁻²sr⁻¹nm⁻ 450 ¹, which is a reasonable error in line with the mission requirements. The results obtained at the O_2 -B for 451

452 the different RT simulations are more similar, and the RMSE is generally lower than 0.05 mWm⁻²sr⁻ 453 1 nm⁻¹ (RRMSEs% < 4) and a slight underestimation is observed.

454

455 <Figure 5>

456

The agreement between the simulated and retrieved SIF and ρ spectra for one of the better result (case 17) and for one of the worst results (case 30) are shown in Figure 6. It can be observed that for the worst case, the true reflectance is not smooth at the O₂-A band but it presents oscillations. The latter are produced by the coupled canopy-atmosphere directional effects when the four reflectance terms from SCOPE are coupled with the atmospheric functions to obtain the total canopy reflectance ρ . This effect produces most of the error in the retrieval of the reflectance spectra and consequently the SIF retrieval results are a bit underestimated at the O₂-A.

464

465 <Figure 6>

466

467 This effect can be also observed by evaluating the different reflectance terms simulated in SCOPE. Figure 7 shows the variability of the four reflectance terms for cases 17 and 30. The four reflectance 468 terms simulated with SCOPE are more similar for case 17 and for all the other cases where better 469 470 retrieval performances were found. On the contrary, the variability observed for case 30 indicates a 471 canopy with a strong anisotropic reflectance, which causes difficulties in the fluorescence estimations. 472 This is explained by the fact that the depth of the atmospheric absorption depends on the path-length 473 followed by photons from the sun via the ground to the sensor. The shortest and most direct photon path through the atmosphere corresponds to the sun-target-sensor route. Therefore the r_{so} corresponds to the 474 475 shortest photon path and the shallowest absorption depth. Reflectance factors involving diffuse fluxes

like the r_{do} , r_{sd} and r_{dd} correspond to the reflection of sky radiation by the target or to hemispherically 476 477 reflected radiation by the surroundings. These are always associated to longer photon paths through the atmosphere, and therefore to deeper absorptions. If the r_{so} term is higher than the other reflectance 478 479 terms, shorter photon paths receive more weight, with less deep atmospheric absorption, which 480 contributes to the in-filling and therefore SIF will be underestimated. If the other reflectance terms are higher than r_{so} , then long photon paths receive more weight, due to deep atmospheric absorption, which 481 482 will work as the negative of infilling. The in-filling by fluorescence can therefore be confused with the 483 apparent positive or negative in-filling caused by directional effects. These effects should be considered 484 in further retrieval algorithms because they have a large impact on the overall accuracy. For example, 485 the spectral fitting approach can be improved considering both direct and diffuse canopy reflectance. 486 Alternatively, Verhoef et al., (2014) proposes the numerical inversion of a simplified version of SAIL 487 canopy RT model to represent the different direct and diffuse reflectance terms.

488

489 <Figure 7>

490

491 A typical example of the reflectance and fluorescence retrieval at the O_2 -B band is shown in Figure 8. 492 Although fluorescence and reflectance have a more complex behaviors in the spectral window 493 considered, the Voigt and the piecewise cubic spline functions are able to fit accurately the two spectral 494 components. The directional effects seem to be less significant in this region probably due to the lower 495 magnitude of oxygen in this absorption band.

496

497 <Figure 8>

499 **4.2 Full Fluorescence spectrum**

The *SpecFit* approach is supposed to provide the full fluorescence emission spectrum by combining the FLORIS-NBS and –WBS radiance observations. The number of knots to be used in the piecewise cubic spline to accurately fit the surface reflectance signatures has been initially evaluated by considering an increased number of knots (Figure 9). The results show that 15 knots can properly model the reflectance in the broad spectral region considered in *SpecFit*. In those cases of reflectance signatures with low chlorophyll content or LAI values (i.e., characterized by a weaker red-edge transition), 12 knots seem to be enough, but in order to assure higher accuracies for all the considered cases 15 knots are preferred.

507

508 <Figure 9>

509

510 The RMSE and RRMSE at the O₂-A, O₂-B and for the entire full SIF spectrum are shown in Figure 10. The average RMSE is generally lower than 0.15 mW m⁻² sr⁻¹nm⁻¹ (RRMSE 10%), and slightly better 511 512 performances are found for cases 15-18 and 31. These results are somewhat close to those found at the 513 O₂-A band by the SFM approach in narrower spectral windows. In fact, the overall accuracy is similar 514 (at least only slightly worse due to the broader spectral window), and more accurate results are found 515 for cases characterized by weaker directional effects. Therefore, the results suggest that the surface-516 atmosphere directional effects at the O_2 -A band also have a large impact on the retrieval accuracy of 517 *SpecFit* algorithm.

518

519 <Figure 10>

520

521 In spite of the difficulties in modeling ρ and SIF in a broader spectral window, the retrieved spectra 522 well represent the target (i.e. input from SCOPE). The full SIF spectrum retrieval results for one of the better cases (case 17) and for the worst (case 30) are reported Figure 11. The reflectance is well modelled in the entire spectrum and it is close to the target spectrum. Also the spectral behavior of retrieved fluorescence is similar to the target, but in some cases slight underestimations occurred mostly driven by the directional effects at the O_2 -A band. It must be noted that the different shapes of the fluorescence strictly depend on the different relative contributions of PS_I and PS_{II} to the total fluorescence spectrum and on the leaf chlorophyll content.

529

530 <Figure 11>

531

532 **4.3** Comparison of the retrievals

The scatterplot between target and retrieved SIF values allow a final performance evaluation of the different retrieval algorithms proposed. The two retrieval approaches show an overall good agreement at 760 nm, 687 nm, and for the spectrally integrated spectrum for all 31 cases (Figure 12). The slopes at 760 nm are slightly different from the 1:1 line, and the intercepts have slightly negative values that are due to the general underestimation of SIF at the O_2 -A band caused by the surface-atmosphere directional effects.

539

541

A comprehensive summary of the retrieval performances relative to both retrieval methods is reported in Table 5. The goodness of fit in terms of the adjusted r^2 are 0.97-0.98 for the O₂-A band. The RMSE (RRMSE) 0.037 mW m⁻² sr⁻¹ nm⁻¹ and 0.044 mW m⁻² sr⁻¹ nm⁻¹ (7.4% and 6.2%) respectively for SFM and *SpecFit*. For the O₂-B, the retrieval performances at 687 nm are generally better, the slope parameter is closer to 1.0, the r^2_{adj} 0.99, and the general RMSE (RRMSE) 0.018 mW m⁻² sr⁻¹ nm⁻¹ (3.0%). It must be noted that results obtained at the oxygen absorption bands by *SpecFit* are a bit better than SFM. This is because the applied fitting over a broader spectral window reduces the directional effects, and instrumental noise, which have a higher impact within the strong absorption bands. Finally, the results achieved by *SpecFit* for the spectrally integrated SIF values in the 670-780 nm spectral range show an overall RMSE (RRMSE) of 6.225 mW m⁻² sr⁻¹ (6.4%).

- 552
- 553 <Table 5>
- 554

555 The overall results presented here are analogous with respect to the earlier works based on the spectral fitting methods (M. Mazzoni et al., 2012; M. Meroni et al., 2010) at the O₂ bands. The piecewise cubic 556 557 spline and Voigt spectral functions, employed to predict reflectance and fluorescence respectively, were 558 also identified as best candidates in the previous work by Mazzoni et al., 2012. However, the 559 quantitative comparison is difficult because the radiative transfer simulations used are very different. In 560 fact, the four-way RT forward model used in this work includes a more realistic representation of the 561 surface directional effect which allowed us to isolate directional effects as the major source of error in 562 the retrieval, especially at the O_2 -A band. For such reason, only those cases which characterize low 563 directional effects (cases 15-18, 31) should be used in the comparison with earlier studies. Furthermore, 564 the instrument specifications in terms of spectral resolution, sampling interval, and SNR assumed in 565 previous works were generally better. The comparison with the results achieved from other methods 566 (Frankenberg et al., 2011; Guanter et al., 2010; Guanter et al., 2012, 2014; Joiner et al., 2011) is even 567 more difficult because the retrieval errors reported also include the additional uncertainty associated 568 with atmospheric compensation. The comparison with Zhao et al., (2014) is also difficult because the results reported are in absolute values (i.e., RMSE) which are dependent on the specific dataset of 569 radiative transfer simulation used. 570

571

572 5 Fluorescence derived indices from FLEX mission

A number of products are planned to be delivered from the FLEX mission. The Level 1 products will consist of calibrated and geometrically corrected top-of-atmosphere radiances. The Level 2 products will be the surface fluorescence maps and derived indices for vegetation status monitoring. Level 3 products will regard spatial mosaics (regional, continental and global scale) and temporal composites (monthly, seasonal and annual). The Level 4 will be the higher level products after the assimilation of fluorescence and canopy variables within dynamical vegetation models for the generation of gross primary productivity maps.

The spectral fitting based algorithms proposed in this work form part of the entire processing needed for converting Level 1 to Level 2 products. A synthetic flowchart of the retrieval scheme for FLEX, starting from the top of atmosphere radiance (Level 1) up to the surface reflectance and fluorescence spectrum (Level 2a) is summarized in Figure 13. As mentioned earlier, the atmospheric correction process relies also on the S3 data that provide valuable information to constrain the atmospheric parameter retrieval. The resulting bottom of atmosphere radiance is thus decomposed into its reflected and fluorescence components.

587

588 <Figure 13>

589

Finally, a number of indices can be routinely derived from the full fluorescence spectrum facilitating the exploitation of the spectral information. In addition to fluorescence values at the O_2 bands (SIF_{687nm}; SIF_{760nm}) which are useful for comparing SIF maps derived from FLEX with maps derived by others methods and sensors (usually centered at 687 and 760 nm), a number of other parameters informative of plant's activity and physiological status can be extracted at characteristic wavelengths (Table 6). For example, the SIF radiance at the maximum of the two emission peaks (maxSIF_{red}; maxSIF_{far-red}), the wavelength where the maximum of the emission peaks occurs ($max\lambda_{red}$, $max\lambda_{far-red}$), the radio between the red and far-red emission peaks (ratioSIF), and the spectrally integrated SIF emission (intSIF).

598

599 <Table 6>

600

601 6 Conclusions

602 The development and testing of fluorescence retrieval algorithms suitable for the FLEX mission, and 603 for general application in future space missions, have been investigated in this study. Two major 604 spectral fitting approaches have been considered: i) the retrieval at the O₂ absorption bands; and ii) the 605 retrieval of the full fluorescence emission spectrum by a novel algorithm. The algorithms have been 606 tested on a state-of-the-art dataset of RT simulations resampled according to the current technical 607 specification of the FLORIS space-borne sensor. For the first time, the RT calculations include a more 608 realistic coupling between the canopy and atmospheric directional effects. The latter resulted in one of 609 the major sources of error in the retrieval. In particular, the retrieval at the O₂-A band was severely 610 affected by the directional effects, while the analysis at O₂-B band showed less sensitivity. Therefore, 611 further studies are needed to better understand and quantify the impact of directional effects on the total 612 error budget and the opportunity of applying the SIF retrieval scheme at TOA level (where the 613 directional effects are larger) will be evaluated. Nevertheless, the results presented suggest the 614 possibility of retrieving red SIF with a high level of accuracy from the FLORIS sensor. The use of 615 mathematical functions able to accurately model SIF and p spectral behaviors, combined with the SR 616 and SNR provided by the FLORIS sensor permitted of retrieving red SIF with a high level of accuracy. 617 The performance of the novel *SpecFit* retrieval algorithm is similar (and sometimes better) than those 618 obtained with SFM at the O₂-A and O₂-B bands. The promising results achieved open new perspectives for further investigating the full fluorescence spectrum in relation to plant species and traits analysis, environmental conditions, structural and physiological variables, photosynthetic rates and stress occurrences. The proposed algorithm can be further optimized to retrieve SIF from high-resolution field spectroscopy measurements (Cogliati et al., 2015; Corp et al., 2010; Daumard et al., 2012; Meroni et al., 2008; Rossini et al., 2010) providing ground-based measurements useful to interpret satellite remote sensing observations.

625 However, the results presented here for both the retrieval approaches, cannot be representative of the 626 true total error budget as expected from FLEX (i.e. from TOA radiance to canopy SIF) because an a-627 priori atmospheric correction was used. However, the FLEX End-to-End Simulator under development 628 will enable a complete understanding of the retrieval error (from TOA radiance to surface reflectance 629 and fluorescence), and the mission performance can be realistically investigated. Further instrument-630 related effects, such as smile, band broadening and stray-light, factors known to affect the retrieval 631 accuracy and which are neglected in this study. This is because they will be minimized by sensor 632 design, high-quality optics, and by controlling the cleanliness during the instrument manufacturing 633 phase.

In summary, we have demonstrated the potential to retrieve the full spectrum of fluorescence from hyperspectral observations which can open new applications for better understanding of the terrestrial environment and ecosystem function.

637

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780 FIGURE CAPTIONS

Figure 1: Typical spectra of top of atmosphere radiance as detected by the FLORIS-WBS (gray) and NBS (blue) (stacked for clarity), sun-induced fluorescence (red) and reflectance (green) are shown in the upper plot. Details of the high-resolution FLORIS-NBS spectra at the O_2 -B (left) and O_2 -A (right) are depicted in the bottom plots.

Figure 2: Atmospheric transfer functions used in the forward model for calculating top of atmosphere radiance (left table), the angular brakets represent the spectral convolution to the sensor spectral bands. Typical spectra of the atmospheric transfer functions used in the forward radiative transfer model are shown in the 500-780 nm spectral range (rigth plot).

Figure 3: Range of variations of reflectance (left), fluorescence (middle) and at sensor radiance
radiance (rigth) of the 31 simulated cases.

Figure 4: Average RMSE over the 31 simulated cases at the O_2 -A (left) and O_2 -B (right) bands for the broader spectral window. The colors refer to three major classes of functions: i) polynomial (red); ii) Legendre polynomial (blue); and iii) piecewise cubic splines (green). The different SFM versions (1-18) are on the abscissa. The box center is the median, the edges the 25th and 75th percentiles and the whiskers extend to the most extreme data points. The scale limit settled at 1 mW m⁻² sr⁻¹ nm⁻¹ to highlight results with higher accuracies, while larger RMSE are grouped on top.

Figure 5: SIF retrieval accuracy in terms of RMSE (top), RRMSE% (middle) and SIF_{diff} (bottom) at the O_2 -A (blue) and O_2 -B (red) bands by means of SFM version 18-S. Cases 1-19 and 20-31 consider variation of surface and atmospheric parameters respetively.

Figure 6: Results achieved by using SFM at the O₂-A band. The spectra for one of the better results (case 17), and for one of the worst results (case 30), are shown on the left and right, respectively. The upper plots show the reflectance quantities: ρ is the target reflectance (blue), ρ_{app} is the apparent reflectance (green), and ρ_{ret} the retrieved one (red). The lower plots show the target (blue) and retrieved (red) SIF.

Figure 7: Four reflectance terms r_{so} (blue), r_{do} (green), r_{sd} (red), r_{dd} (gray) in the visible to near-infrared for the database cases 17 (left) and case 30 (rigth).

Figure 8: Spectral fitting retrieval at the O₂-B band for case 17. Left plot shows the target ρ (blue), the apparent ρ_{app} (green) and the retrieved ρ_{ret} (red) reflectance. The right plot shows target (blue) and retrieved (red) SIF spectra.

Figure 9: Optimization of the piecewise cubic spline used for fitting surface reflectance. The left plot shows the RMSE for simulations 1-18 (where surface reflectance has been changed) considering different knots numbers (3-20 knots). The right plot shows a typical fitting of surface reflectance (ρ) by using piecewise cubic spline with 15 knots (ρ_{spline}).

Figure 10: Retrieval accuracies in terms of RMSE (top), RRMSE% (bottom) at the O₂-A (blue), O₂-B (red) and for the full fluorescence spectrum (green) achieved by SpecFit algorithm. The cases 1-19 consider variation of surface parameters, while cases 20-31 variations of atmospheric parameters.

Figure 11: Retrieval of full SIF spectrum with SpecFit algorithm in the 670-780 nm fitting range. The upper plots show target reflectance ρ (blue), apparent reflectance ρ_{app} (green), and retrieved reflectance ρ_{ret} (red line). The charts on the bottom show target (blue) and retrieved (red) SIF spectra. The plots on the left and right refer to case 17 (i.e., better results) and case 30 (i.e., worse result) respectively.

- 821 **Figure 12:** Comparison between SIF values retrieved by spectral fitting at the O₂ bands (blue circles)
- 822 and SpecFit (red squares). The target (SIF) and retrieved (SIF_{ret}) fluorescence values at 760 nm (left),
- 823 687 nm (center) and for the integral of the full SIF emission spectrum (right) obtained on the 31 RT
- simulations are shown. The gray dash-dot represents the 1:1 line.
- Figure 13: Flow chart of the FLEX retrieval scheme from top of atmospfere radiance to the surfacereflectance, fluorescence, and derived indices.

Table 1: Technical characteristics of the FLORIS spectra in terms of spectral resolution (SR), spectral sampling interval (SSI), and signal to noise ratio (SNR) for the different spectral regions.

Spectral Region	Visible	SI	Fred	ree	red-edge		SIF _{far-red}			
λ (nm)	500-677	677-686	686-697	697-740	740-755	755-759	759-762	762-769	769-780	
SR (nm)	3.0	0.6	0.3	2.0	0.7			0.3	0.7	
SSI (nm)	2.0	0.5	0.1	0.65	0.5		0.1		0.5	
SNR	245	340	175	425	linear from 510 to 1015	1015	115	linear from 115 to 455	1015	

Table 2: Surface and atmospheric parameters for SCOPE and MODTRAN radiative transfer models. The low, standard and high values used in the different cases simulated are reported. The right column "Coupling" indicates whether the parameters for the two models are coupled.

				Values		
Case	Parameter	Unit	low	medium (case 19)	high	Coupling
Surface par	rameters (SCOPE)					
1-2	Soil reflectance		1	2	3	Ν
3-4	chlorophyll content (Cab)	μg cm ⁻²	20	40	80	N
5-6	leaf water equivalent layer (Cw)	cm	0.01	0.02	0.03	Ν
7-8	dry matter content (Cdm)	g cm ⁻²	0.0025	0.005	0.01	Ν
9-10	senescent material content (Cs)	-	0.05	0.10	0.20	N
11-12	maximum carboxylation capacity (Vcmax)	µmol m ⁻² s ⁻¹	0	40	100	N
13-14	stomatal conductance (m)	-	2	5	9	Ν
15-16	leaf area index (LAI)		0.5	2	8	Ν
17-18	leaf inclination distribution function (LIDF)		planophile	spherical	erectophile	N
Atmospher	ic parameters (MODTRA	AN5)				
20-21	Surface height	m	0	400	1200	Y
22-23	Visibility	km	5	20	80	Ν
24-25	Humidity		0.5 imes	$1.0 \times$	2.0×	Y
26-27	Aerosol Type		maritime	rural	urban	N
28-29	Profile		mid-latitude winter	mid-latitude summer	tropical	Y
30-31	Solar zenith angle	deg	30	45	60	Y

Table 3: Spectral windows at O₂-B and O₂-A bands tested with different spectral fitting algorithms.

spectral	O ₂ - B	O ₂ -A
window	[nm]	[nm]
1	686-691	759-769
2	686-696	750-780

Table 4: Combinations of mathematical functions tested in the SFM algorithms to retrieve SIF and ρ .

Polynomial			Legendre			Spline		
ID	ρ	SIF	ID	ρ	SIF	ID	ρ	SIF
1-P	Linear	Linear	1-L	Linear	Linear	1-S	Linear	Linear
2-P	Quadratic	Linear	2-L	Quadratic	Linear	2-S	Quadratic	Linear
3-P	Cubic	Linear	3-L	Cubic	Linear	3-S	Cubic	Linear
4-P	Linear	Quadratic	4-L	Linear	Quadratic	4-S	Linear	Quadratic
5-P	Quadratic	Quadratic	5-L	Quadratic	Quadratic	5-S	Quadratic	Quadratic
6-P	Cubic	Quadratic	6-L	Cubic	Quadratic	6-S	Cubic	Quadratic
7-P	Linear	Cubic	7-L	Linear	Cubic	7-S	Linear	Cubic
8-P	Quadratic	Cubic	8-L	Quadratic	Cubic	8-S	Quadratic	Cubic
9-P	Cubic	Cubic	9-L	Cubic	Cubic	9-S	Cubic	Cubic
10-P	Linear	Gaussian	10-L	Linear	Gaussian	10-S	Linear	Gaussian
11-P	Quadratic	Gaussian	11-L	Quadratic	Gaussian	11-S	Quadratic	Gaussian
12-P	Cubic	Gaussian	12-L	Cubic	Gaussian	12-S	Cubic	Gaussian
13-P	Linear	Lorentzian	13-L	Linear	Lorentzian	1 3- S	Linear	Lorentzian
14-P	Quadratic	Lorentzian	14-L	Quadratic	Lorentzian	14-S	Quadratic	Lorentzian
15-P	Cubic	Lorentzian	15-L	Cubic	Lorentzian	15-S	Cubic	Lorentzian
16-P	Linear	Voigt	16-L	Linear	Voigt	16-S	Linear	Voigt
17-P	Quadratic	Voigt	1 7- L	Quadratic	Voigt	17-S	Quadratic	Voigt
18-P	Cubic	Voigt	18-L	Cubic	Voigt	18-S	Cubic	Voigt

Table 5: Comparison between SFM and *SpecFit* algorithms at 760 nm, 687 nm and for the integral of the full fluorescence emission spectrum (670-780 nm) for the 31 cases. The retrieval accuracy in terms of: slope (c_1) and intercept (c_2) of the linear model, r^2 adjusted, RMSE ([mW m⁻²sr⁻¹nm⁻¹] and [mW m⁻²sr⁻¹] for values at the O₂ bands and spectral integral respectively), and RRMSE% are reported.

	SFM _{760nm}	SFM _{687nm}	SpecFit _{760nm}	SpecFit _{687nm}	SpecFit _{Int}
<i>c</i> ₁	0.935*	0.978*	0.956*	0.956*	0.936*
<i>c</i> ₂	-0.007	0.035	-0.012	0.081*	1.663
r ² adj	0.98	0.99	0.97	0.99	0.97
RMSE	0.037	0.018	0.044	0.018	6.225
RRMSE%	7.4	3.0	6.2	2.9	6.4

* p-values < 0.05

Table 6: List of indices derived from the full fluorescence spectrum.

Derived products	Unit	Specification
SIF _{760nm}	$mW m^{-2}sr^{-1}nm^{-1}$	Fluorescence at 760 nm
SIF _{687nm}	$mW m^{-2}sr^{-1}nm^{-1}$	Fluorescence at 687 nm
<i>max</i> SIF _{far-red}	$mW m^{-2} sr^{-1} nm^{-1}$	Fluorescence at the maximum of the far-red peak
maxSIF _{red}	$mW m^{-2}sr^{-1}nm^{-1}$	Fluorescence at the maximum of the red peak
$max\lambda_{far-red}$	nm	Wavelength of the maximum for far-red peak
$max\lambda_{red}$	nm	Wavelength of the maximum for red peak
ratioSIF	-	maxSIF _{red} / maxSIF _{far-red}
IntSIF	$mW m^{-2} sr^{-1}$	Integral of the full SIF emission



Atmospheric Functions	Name
$\langle E_s^o \rangle \cos \theta_s / \pi$	t ₁
$\langle \rho_{so} \rangle$	t2
$\langle \rho_{dd} \rangle$	t3
(τ _{ss})	t4
$\langle \tau_{sd} \rangle$	£5
(T ₀₀)	t ₆
$\langle \tau_{do} \rangle$	£7
$(\tau_{ss}\tau_{oo})$	t _{si}
$\langle \tau_{sd} \tau_{oo} \rangle$	t9
$(\tau_{ss}\tau_{do})$	t ₁₀
$\langle \tau_{sd} \tau_{do} \rangle$	£11
$\langle \tau_{ss} \rho_{dd} \rangle$	t ₁₂
$\langle \rho_{dd} \tau_{oo} \rangle$	t ₁₃
$\langle \tau_{ss} \rho_{dd} \tau_{oo} \rangle$	t14







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Figure 6 (size 1.5 -column) Click here to download high resolution image













Figure 11 (size 1.5-column) Click here to download high resolution image





